## <u>Spherically Symmetric</u> <u>Charge Densities</u>

Consider volume charge densities  $\rho_{\nu}(\bar{r})$  that are functions of spherical coordinate r only, e.g.:

 $\rho_{\nu}(\overline{\mathbf{r}}) = \frac{1}{r^{2}} \quad \text{or} \quad \rho_{\nu}(\overline{\mathbf{r}}) = e^{-r}$ 

We call these types of charge densities **spherically symmetric**, as the charge density changes as a function of the distance from the origin only (i.e., is independent of coordinates  $\theta$  or  $\phi$ ).

As a result, the charge distribution in this case looks sort of like a "**fuzzy ball**", centered at the origin!

Using the point form of Gauss's Law, we find the resulting static electric field **must** have the form:

 $\mathbf{E}(\overline{\mathbf{r}}) = \mathcal{E}(\mathbf{r}) \hat{a}_{r}$  (for spherically symmetric  $\rho_{\nu}(\overline{\mathbf{r}})$ )

Think about what this says. It states that the resulting static electric field from a spherically symmetric charge density is:

\* A function of spherical coordinate ronly.

\* Points in the direction  $\hat{a}_r$  (i.e., away from the origin at every point).

As a result, we can use the integral form of Gauss's Law to determine the specific scalar function E(r) resulting from some specific, spherically symmetric charge density  $\rho_v(\bar{r})$ .

Recall the integral form of Gauss's Law:

Consider now a surface 5 that is a **sphere** with radius *r*, centered at the origin. We call this surface the **Gaussian Surface** for spherically symmetric charge densities.

To we why, we integrate over this Gaussian surface and find:

$$\bigoplus_{s} \mathbf{E}(\overline{\mathbf{r}}) \cdot \overline{ds} = \int_{0}^{2\pi\pi} \int_{0}^{\pi} \mathbf{E}(\overline{\mathbf{r}}) \cdot \hat{a}_{r} r^{2} \sin\theta d\theta d\phi$$

$$= \int_{0}^{2\pi\pi} \int_{0}^{\pi} \mathbf{E}(r) \hat{a}_{r} \cdot \hat{a}_{r} r^{2} \sin\theta d\theta d\phi$$

$$= E(r) r^{2} \int_{0}^{2\pi\pi} \int_{0}^{\pi} \sin\theta d\theta d\phi$$

$$= 4\pi r^{2} E(r)$$

$$4\pi r^{2} E(r) = \frac{Q_{enc}}{\varepsilon_{0}}$$

Rearranging, we find that the function E(r) is:

$$E(r) = \frac{Q_{enc}}{4\pi\varepsilon_0 r^2}$$

The enclosed charge  $Q_{enc}$  can be determined for a spherically symmetric distribution (a function of r only!) as:

$$Q_{enc} = \iiint_{V} \rho_{v}(\overline{\mathbf{r}}) dv$$
$$= \int_{0}^{2\pi} \int_{0}^{\pi} \int_{0}^{r} \rho_{v}(r') r'^{2} \sin \theta dr' d\theta d\phi$$
$$= 4\pi \int_{0}^{r} \rho_{v}(r') r'^{2} dr'$$

Therefore, we find that the static electric field produced by a **spherically symmetric** charge density is  $\mathbf{E}(\overline{\mathbf{r}}) = E(\mathbf{r})\hat{a}_r$ , where the scalar function  $\mathbf{E}(\mathbf{r})$  is:

$$E(r) = \frac{Q_{enc}}{4\pi\varepsilon_0 r^2}$$
$$= \frac{1}{\varepsilon_0 r^2} \int_0^r \rho_v(r') r'^2 dr'$$

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Or, more specifically, we find that the static electric field produced by some **spherically symmetric** charge density  $\rho_{\nu}(\overline{r})$  is:

 $\mathbf{E}(\mathbf{\bar{r}}) = \frac{Q_{enc}}{4\pi\varepsilon_0 r^2} \hat{a}_r$  $= \frac{\hat{a}_r}{\varepsilon_0 r^2} \int_0^r \rho_v(\mathbf{r}') \mathbf{r}'^2 d\mathbf{r}'$ 

Thus, for a **spherically symmetric** charge density, we can find the resulting electric field **without** the difficult integration and evaluation required by **Coulomb's Law**!